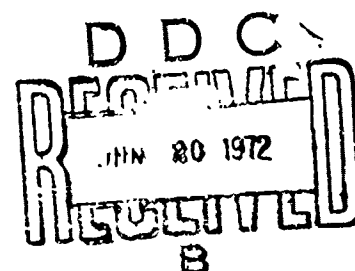


**DYNAMIC PROPERTIES OF POLY (METHYLMETHACRYLATE)
(PMMA) (PLEXIGLAS)**

by

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INTRODUCTION

This study of poly(methylmethacrylate) (PMMA) was conducted under the PREDIX program initiated by the Defense Nuclear Agency.⁽¹⁾ The program was concerned primarily with the response of metals under impulsive loading conditions but some tests were performed on PMMA to provide data on a viscoelastic, amorphous material. Although the tests were not as extensive as those conducted on 6061-T6 aluminum,⁽²⁾ alpha titanium,⁽³⁾ OFHC copper,⁽⁴⁾ and tantalum,⁽⁵⁾ sufficient data were obtained in various areas of material response to warrant preparation of a separate report.

The primary areas studied were: (1) uniaxial stress behavior; (2) elastic constants; (3) equation of state; (4) shock wave profiles; and (5) spall fracture. The experimental techniques utilized in this study are reviewed in Reference 6. Briefly, the principal techniques applied were:

Stress-Strain Studies - A laboratory-type universal testing machine was used for rates $< 0.1/\text{sec}$, a medium strain rate machine for rates of $10^{-3}/\text{sec}$ to $10^2/\text{sec}$, and a Hopkinson bar device for rates of about 10^2 to $5 \times 10^3/\text{sec}$.

Elastic Constants - The pulse superposition method was used for measuring longitudinal and shear wave velocities as a function of temperature.

Equation of State - Gun-launched, flat-plate impact techniques were used to generate uniaxial strain conditions.

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Velocities up to 0.6 mm/ μ sec were achieved with 63.5 mm and 102 mm single-stage compressed-gas guns. Hugoniot data were obtained with x-cut quartz gages.

Wave Propagation and Spall Fracture - Flat-plate impact techniques were also used. Wave profiles were measured with x cut quartz gages, laser velocity interferometer and magnetic foil gages.

MATERIAL PROPERTIES

Poly(methylmethacrylate) or PMMA is a cast thermoplastic acrylic polymer made from methyl methacrylate monomers. The material used for the tests reported here was Rohm and Haas Plexiglas II UVA, an ultraviolet absorbing grade.⁽⁷⁾ The material was purchased as sheet stock in thicknesses of 1.6, 3.2, 12.7 and 25.4 mm. Unless noted, all data reported are for a test direction normal to the plane of the sheet.

The average measured density was 1.183 g/cc (0.0427 lb/in.³). The material has a softening or heat-distortion temperature in the range of 95-105°C and a Rockwell hardness of M93.⁽⁷⁾ Thermal analyses by Asay, et al.,⁽⁸⁾ on cast Acrylux (Cadillac Plastic and Chemical Co.) rod gave the following:

Volume coefficient of thermal expansion, β :

$$145 \times 10^{-6} + 1.4 \times 10^{-6}T \quad /^{\circ}\text{C}$$

(10 to 100°C)

Specific Heat, c_p :

$$0.1974 + 0.00176T \quad \text{cal/g-}^{\circ}\text{C}$$

(-100 to 45°C)

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The above give $\beta = 163 \times 10^{-6}$ and $c_p = 0.23$ at 20°C . Other data on PMMA give values of β as high as 240×10^{-6} and c_p as high as 0.35. ^(7,9)

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SECTION I

UNIAXIAL STRESS BEHAVIOR

The uniaxial stress response of PMMA has been studied under several programs at General Motors Corporation. Results of these studies are given in References 10 to 14, and cover such areas as strain rate and temperature effects in compression and tension, fracture, specimen size effects, and the influence of prestraining (uniaxial stress) and preshocking (uniaxial strain) on stress-strain behavior. In addition, numerous papers have been published on deformation characteristics of PMMA (see, e.g., References 15 and 16).

The response of PMMA to uniaxial stress loading is viscoelastic with high strain rate sensitivity, as shown in Figure 1.* The low-strain modulus and the high-strain flow stress increase by a factor of two or more for a strain-rate increase of 0.003/sec to 1500/sec, and decrease correspondingly for a temperature increase of 22°C to 82°C.

PMMA exhibits significant stress relaxation in uniaxial stress, and a specimen loaded in compression to $\sim 1\%$ strain at $\sim 10/\text{sec}$ relaxed from 2.34 kbar to 1.63 kbar in ~ 0.025 sec.⁽¹¹⁾ Specimens loaded in compression at 2.5/sec to 99% of maximum stress and then unloaded showed no change in stress-strain response on subsequent reloadings.⁽¹¹⁾

* The data in Figure 1 are for 9.5 mm dia. by 12.7 mm long cylinders machined from 12.7 mm dia. cast rod.

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Uniaxial stress tests on unshocked (as-received) and shocked (uniaxial strain compressive loading to 5 kbar) specimens showed no differences in compressive stress-strain response for strain rates of 0.003/sec. to 1160/sec. (11)

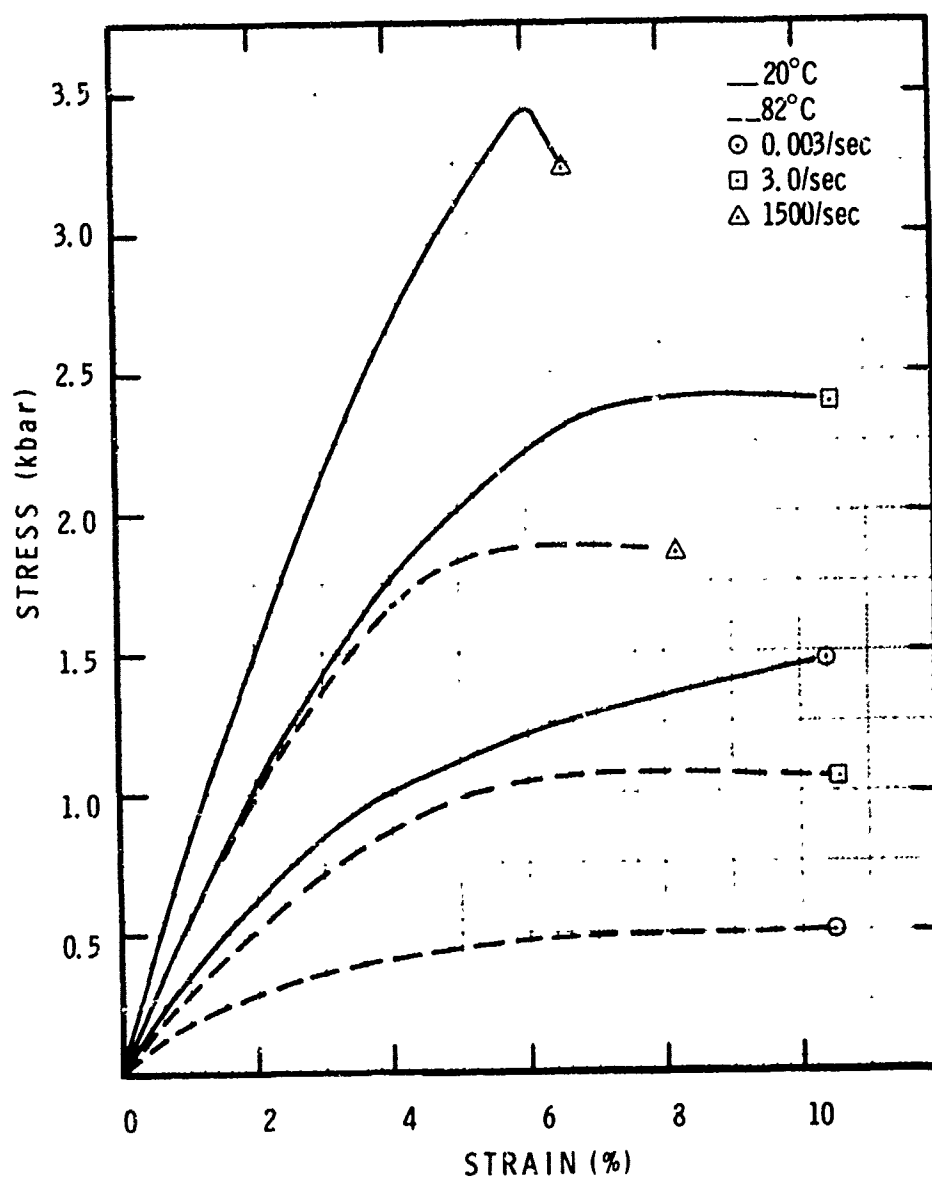


Figure 1 Compressive Stress vs. Strain, PMMA

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SECTION II

ELASTIC BEHAVIOR

WAVE VELOCITIES

Measurements were made of longitudinal and shear wave velocities in PMMA at 20°C and atmospheric pressure ($P = 0$). Values obtained were:

$$\left. \begin{aligned} C_L &= 2.746 \pm 0.003 \text{ mm}/\mu\text{sec} \\ C_S &= 1.392 \pm 0.002 \text{ mm}/\mu\text{sec} \end{aligned} \right\} \quad (1)$$

The temperature dependence of C_L and C_S was also measured and, after correction for thermal expansion, gave:

$$\left. \begin{aligned} C_L &= 2.801 - 26.4 \times 10^{-4}T - 0.05 \times 10^{-4}T^2 \\ C_S &= 1.420 - 13.6 \times 10^{-4}T - 0.028 \times 10^{-4}T^2 \end{aligned} \right\} \quad (2)$$

(T in °C)

Asay, et al., ^(8,17) measured the temperature dependence of C_L and C_S along the axial direction of cast PMMA rod and obtained values in reasonable agreement with those given above.

Asay, et al., ⁽¹⁷⁾ also reported the pressure dependence of C_L and C_S giving:

$$\left. \begin{aligned} C_L &= 2.7519 + 0.2452P - 0.0166P^2 + 0.000589P^3 \\ C_S &= 1.3977 + 0.0965P - 0.00705P^2 + 0.00270P^3 \end{aligned} \right\} \quad (3)$$

(P in kbar)

Note that the zero pressure values of C_L and C_S (2.7519 and 1.3977) are about $\sim 0.3\%$ higher than those given in Equation 1 for material taken from sheet stock.

ELASTIC CONSTANTS

Assuming PMMA is homogeneous and isotropic, elastic wave velocity data can be used to calculate various elastic constants. At 20°C and zero pressure, the following adiabatic constants were obtained:

Bulk wave velocity, $C_B = 2.23 \text{ mm}/\mu\text{sec}$

Sound wave velocity, $C_E = 2.27 \text{ mm}/\mu\text{sec}$

Poisson's ratio, $\nu = 0.327$

Bulk modulus, $K = 58.6 \text{ kbar}$

Shear modulus, $G = 22.9 \text{ kbar}$

Elastic modulus, $E = 60.8 \text{ kbar}$

Lamé's parameter, $\lambda = 43.4 \text{ kbar}$

Isothermal values of K and G as well as adiabatic and isothermal pressure and temperature derivatives at 20°C and zero pressure were calculated using pressure dependence, specific heat and thermal expansion data* reported by Asay, et al.^(8,17) Results are given in Table I.

* Use of higher values for β and c_p than reported by Asay has a small effect on the calculations, with the results being 0.2 to 4% different for $\beta = 240 \times 10^{-6}$ and $c_p = 0.35$.

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TABLE I
ELASTIC CONSTANTS FOR PMMA
(P = 0, T = 20°C)

PARAMETER	VALUE
K^S	58.6 kbar
$(\frac{\partial K^S}{\partial P})_T = K_{0T}^{S'}$	12.7
$(\frac{\partial K^S}{\partial T})_P$	- 0.129 kbar/°C
$(\frac{\partial K^S}{\partial P})_S = K_{0S}^{S'}$	12.2
$(\frac{\partial K^S}{\partial T})_S$	2.85 kbar/°C
K^T	56.4 kbar
$(\frac{\partial K^T}{\partial P})_T = K_{0T}^{T'}$	13.0
$(\frac{\partial K^T}{\partial T})_P$	- 0.151 kbar/°C
$(\frac{\partial K^T}{\partial P})_S = K_{0S}^{T'}$	12.4
$(\frac{\partial K^T}{\partial T})_S$	2.90 kbar/°C
$G^S = G^T$	22.9 kbar
$(\frac{\partial G}{\partial P})_T = G_{0T}'$	3.58
$(\frac{\partial G}{\partial T})_P$	- 0.052 kbar/°C
$(\frac{\partial G}{\partial P})_S = G_{0S}'$	3.36
$(\frac{\partial G}{\partial T})_S$	0.79 kbar/°C

SECTION III

EQUATION OF STATE

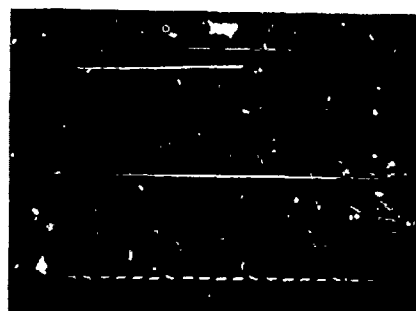
The hugoniot equation of state is the locus of equilibrium states reached after shock-loading of a material, and data are usually obtained either as stress-particle velocity points or as shock velocity-particle velocity points. The hugoniot data presented in this report were obtained with x-cut quartz gages and representative records for direct impact and transmitted wave tests are shown in Figure 2. The buffered direct impact method (tungsten carbide or 4340 steel buffer plate on front of the quartz) permitted stresses up to 20 kbar in PMMA while keeping impact velocity requirements within the capability of the 102 mm gun. The use of quartz gages for transmitted wave tests was primarily for the study of compressive wave development, and results are discussed in the Wave Propagation Section.

The hugoniot may be expressed in several forms and the most convenient form based on quartz gage results is that established by a least-squares fit to data in the stress-particle velocity ($\sigma_H - u_p$) plane. Results for PMMA are given in Figure 3 and the hugoniot is

$$\sigma_H = 29 u_p + 80.3 u_p^2 - 244.5 u_p^3 + 256.4 u_p^4 \quad (4)$$

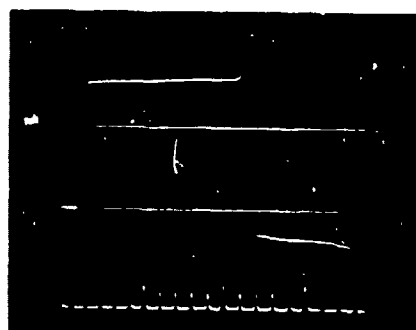
For metals, a second order equation is usually sufficient to give a reasonable fit to the data,⁽²⁻⁵⁾ however, a fourth-order equation is given for PMMA for two reasons. First, the standard error for second and third-order fits is 0.1 kbar, while the fourth-order equation reduces this to 0.05 kbar. Second, the work of other investigators has established that

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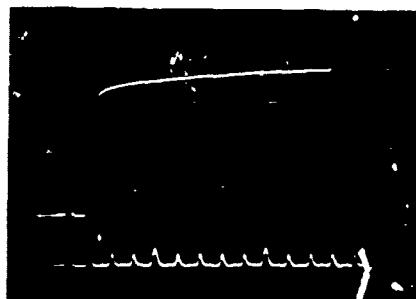


DIRECT IMPACT
PMMA → Qz
 $V_I = 0.122 \text{ mm}/\mu\text{s}$
 $\sigma_H = 3.40 \text{ kbar}$

TIME →



DIRECT IMPACT, BUFFERED
PMMA → 4340 ST./Qz
 $V_I = 0.163 \text{ mm}/\mu\text{s}$
 $\sigma_H = 5.55 \text{ kbar}$



TRANSMITTED WAVE
6061-T6 Al → PMMA/Qz
 $V_I = 0.252 \text{ mm}/\mu\text{s}$
 $\sigma_H = 7.5 \text{ kbar}$

Figure 2 Quartz Gage Records

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PMMA does not have a conventional hugoniot in the low-pressure region (see e.g., Refs. 18-22).^{*} Also shown in Figure 3 are data of Barker and Hollenbach⁽²¹⁾ which show good agreement with the present work.

Transformation of the hugoniot into other planes, such as stress-volume ($\sigma_H - v$) or shock velocity-particle velocity ($U_S - u_p$) is performed by assuming a material model. Since PMMA does not show an elastic-plastic structure, this was done by assuming an ideal single-wave structure with equilibrium initial and final states and applying the mass and momentum conservation equations:

$$\sigma_H = \rho_0 U_S u_p$$

and

$$v = v_0 \left(1 - \frac{u_p}{U_S} \right)$$

Results are given in Table II and Figures 4 and 5. Barker's data and fits⁽²¹⁾ are also shown. The calculated shock velocities from the present work are slightly lower than Barker's measured values, but the agreement is good considering the assumption underlying the data transformation and the effects of material variability and experimental error. Shock velocities measured on three transmitted wave tests are also given in Figure 5.

* Deal⁽²³⁾ did a review paper on all the data available up to 1965 on PMMA-type materials, but the data scatter did not permit resolution of details of low-pressure response.

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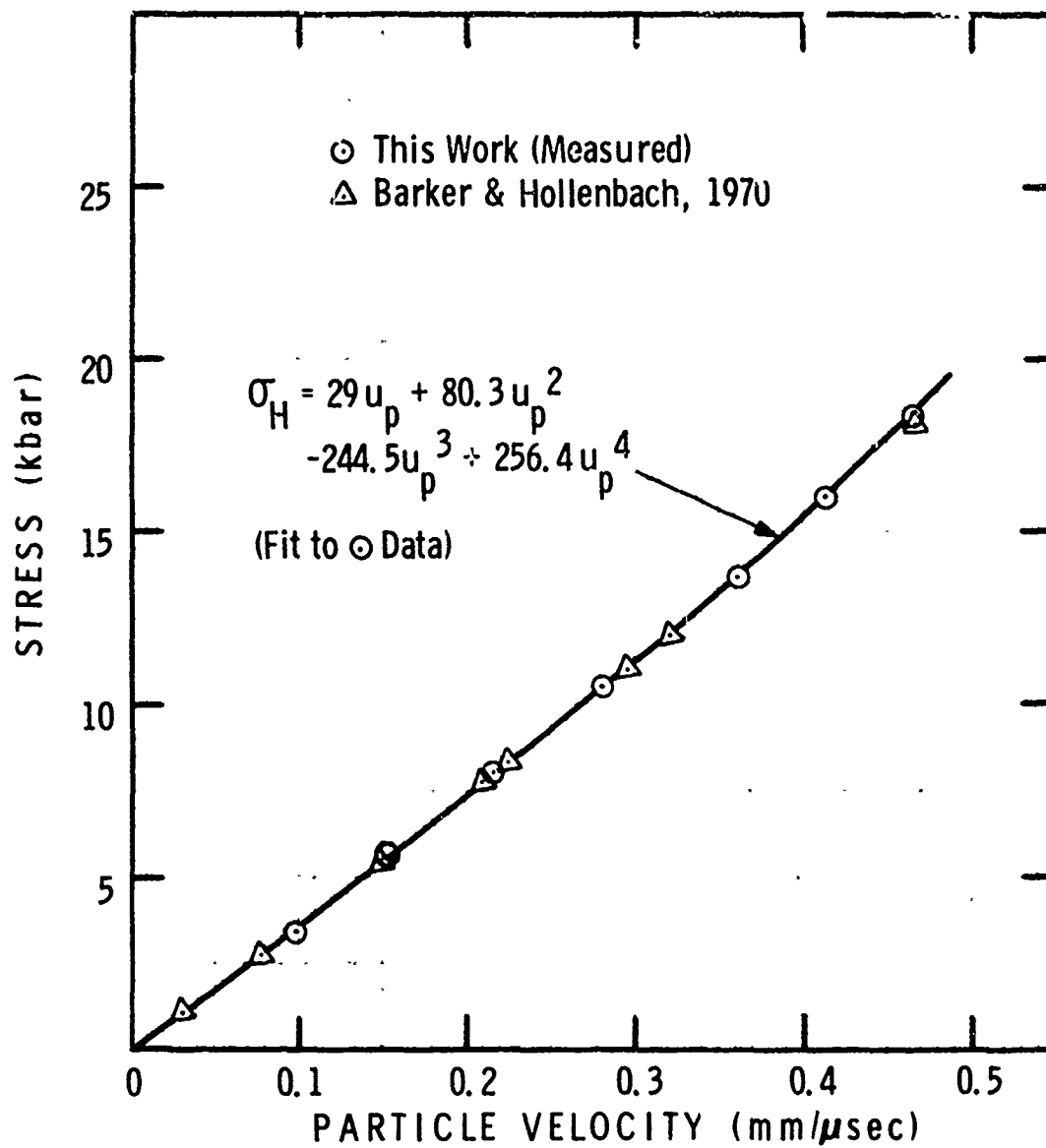


Figure 3 Stress-Particle Velocity Hugoniot, PMMA

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TABLE II
HUGONIOT DATA, PMMA

Stress ¹ (kbar)	Particle Velocity ¹ (mm/ μ s)	Shock Velocity ² (mm/ μ s)	Strain ² $1 - (v/v_0)$	Compression ² $(v_0/v) - 1$
3.40	0.099	2.90	0.0341	0.0353
5.55	0.151	3.11	0.0486	0.0511
8.04	0.214	3.18	0.0674	0.0723
10.6	0.280	3.20	0.0875	0.0959
13.7	0.359	3.23	0.111	0.125
16.0	0.414	3.27	0.127	0.145
18.4	0.468	3.32	0.141	0.164

1. Measured (quartz gages).
2. Calculated ($\rho_0 = 1.183 \text{ g/cc}$).

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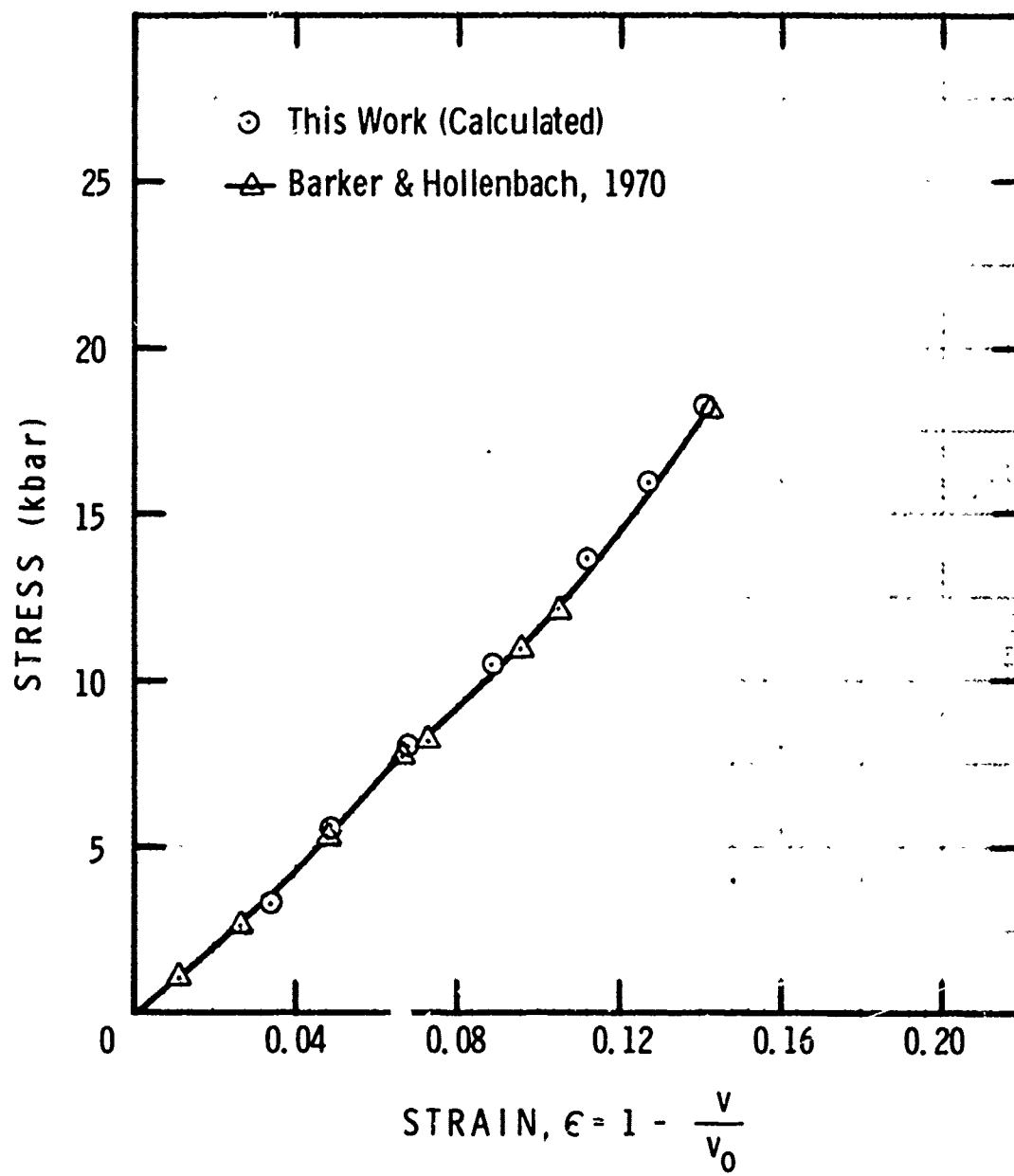


Figure 4 Stress-Strain Hugoniot, PMMA

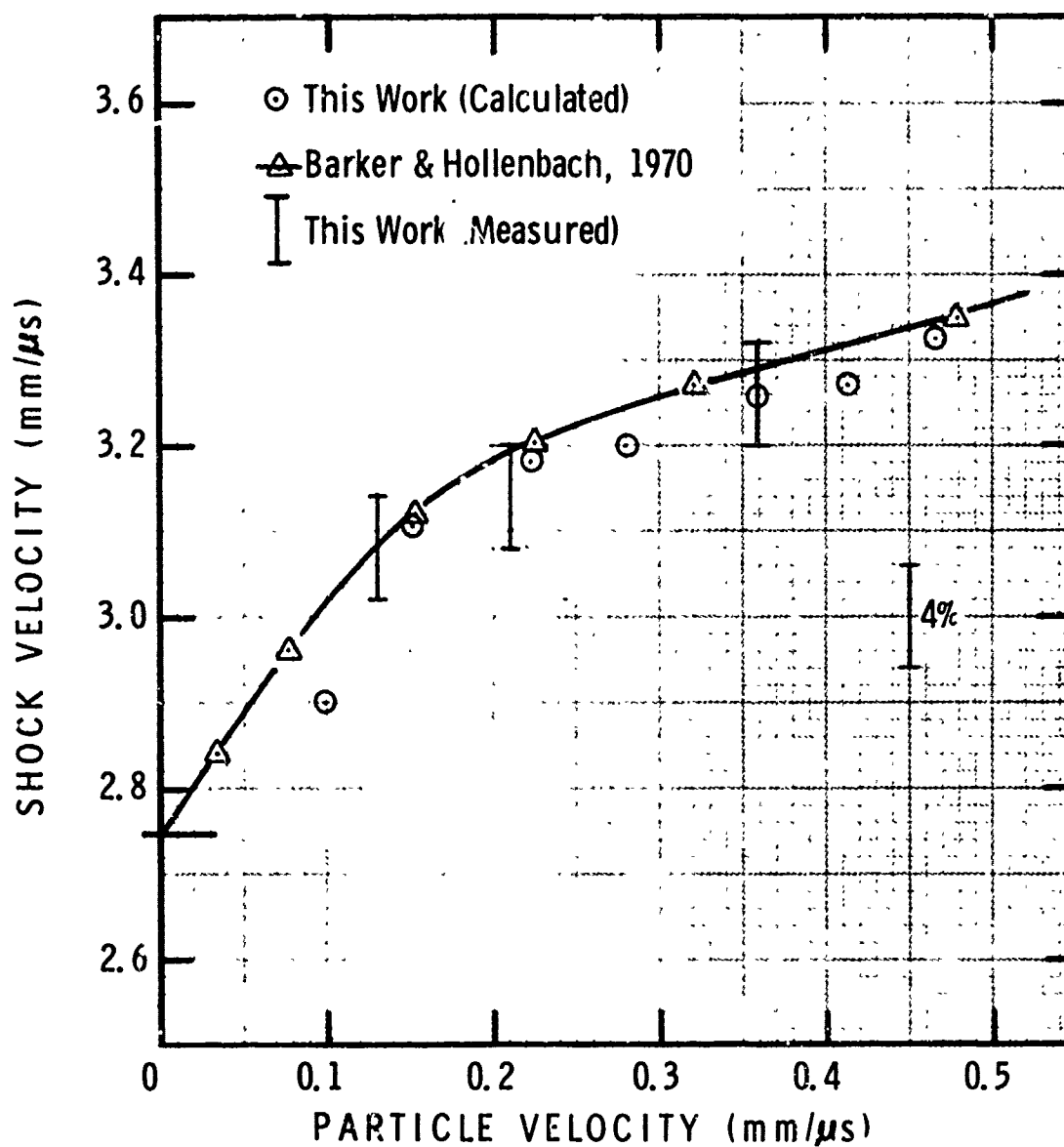


Figure 5 Shock Velocity-Particle Velocity Hugoniot, PMMA

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The stress-compression hugoniot data are compared to ultrasonic hydrostats and hydrostatic compressibility in Figure 6. The ultrasonic hydrostats are based on Asay's pressure data and the equations of Birch, (24,25)

$$P_T = 84.6 (1+\mu)^{7/3} - (1+\mu)^{5/3} + 6.75 (1+\mu)^{2/3} - 1 \quad (5)$$

Murnaghan, (26)

$$P_S = 4.80 (1+\mu)^{12.2} - 1 \quad (6)$$

and Keane. (27,28)

$$P_S = 7.9 \left[(1+\mu)^{9.5} - 1 \right] - 16.7 \ln (1+\mu) \quad (7)$$

The hydrostatic compression curve is Bridgman's data⁽²⁹⁾ as reported by Halpin and Graham.⁽¹⁹⁾ Corrections from isothermal (P_T) and adiabatic (P_S) conditions to the hydrostat are less than 1% at 20 kbar and have been neglected. Compared to metals,⁽²⁻⁵⁾ the agreement between the shock wave data and the hydrostats is relatively poor in the low-pressure region. The ultrasonic and hydrostatic compression results, which are essentially "static" measurements, do not show the inflections in stress-compression behavior indicated by the dynamic shock wave measurements.

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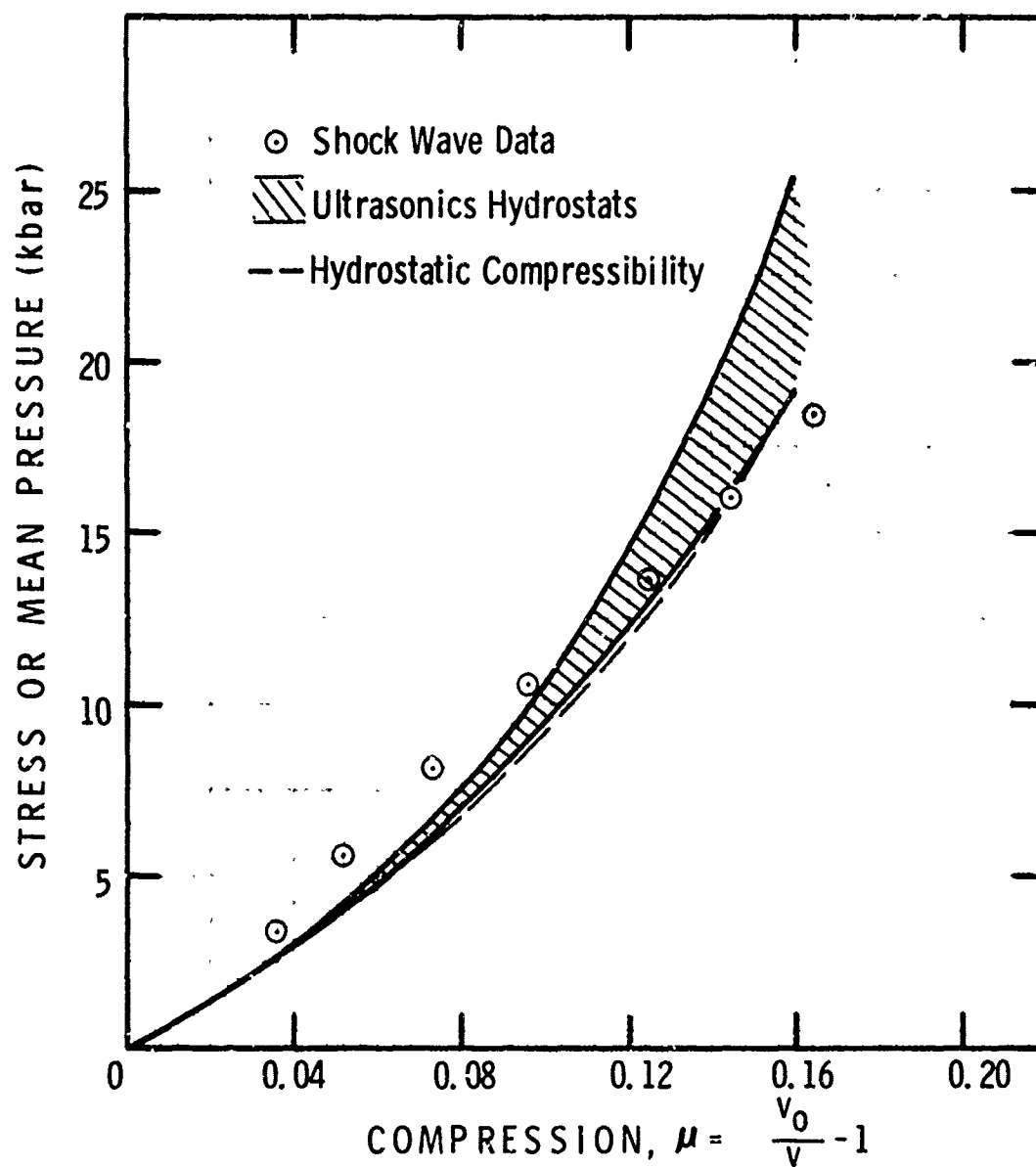


Figure 6 Stress-Compression Hugoniot, PMMA

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SECTION IV

WAVE PROPAGATION

Shock wave characteristics were measured using quartz, a velocity interferometer and magnetic foil gages. The results presented here complement the more detailed work reported by Barker and Hollenbach⁽²¹⁾ and Schuler⁽²²⁾ on the same grade of PMMA. They studied compressive and release wave characteristics at stresses up to 25 kbar and propagation distances out to 37 mm and also examined steady-state wave conditions.

Quartz gage records of compressive waves in PMMA at propagation distances of 3 and 6 mm are shown in Figure 7. The rise-time (after correcting for tilt effects) is only a few nanoseconds up to about two-thirds peak stress, after which there is a gradual rounding of the wave front for $\sim 0.5 \mu\text{sec}$.

Representative records from the velocity interferometer and magnetic foil gage tests are given in Figure 8. Interferometer fringe patterns were generally clean and well-defined, giving good release wave arrival time and shape, although the compressive wave rise-time was too fast to permit fringe resolution. The magnetic foil gage test records were also noise-free but did show a slight "hump" at the first release level which has not been explained.

Velocity interferometer records are shown in Figure 9 with test conditions for each shot listed in Table III. The compressive behavior is essentially the same as described above although there is less rounding for the low velocity test (No. 95, ~ 1.7 kbar). Spreading of the release wave with

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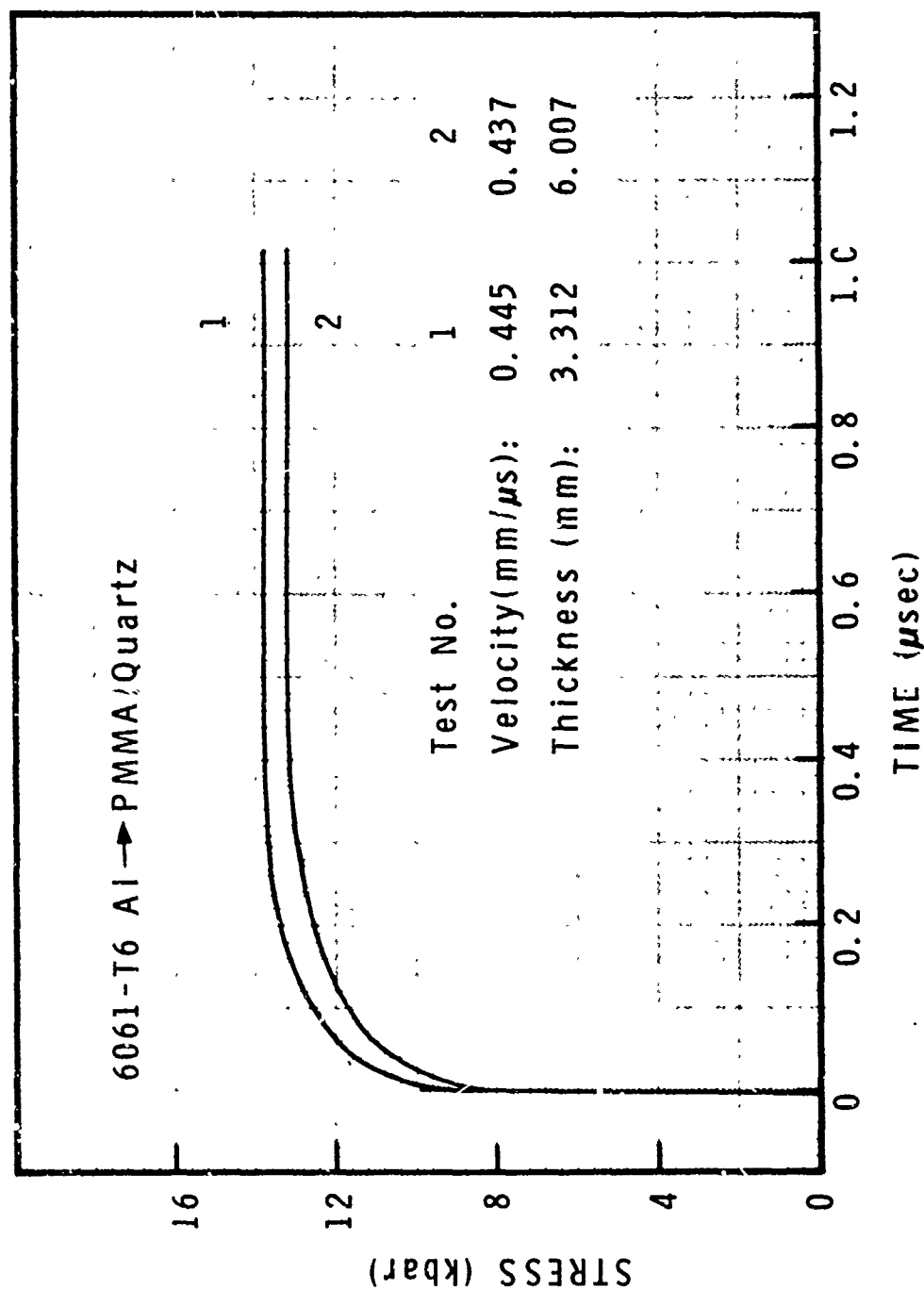
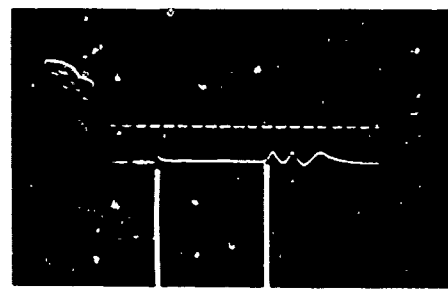


Figure 7 Compressive Waves in PMMA (X-Cut Quartz Gages)

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VELOCITY INTERFEROMETER



PMMA → PMMA/PMMA

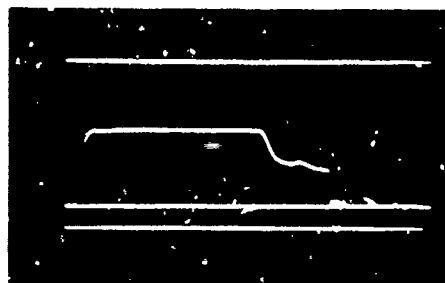
$V_1 = 0.407 \text{ mm}/\mu\text{s}$

$X = 3.414 \text{ mm}$

$\sigma_H = 7.5 \text{ kbar}$

Release Arrival
Compressive Wave

MAGNETIC FOIL GAGE



FQ → PMMA/M. F. /PMMA

$V_1 = 0.163 \text{ mm}/\mu\text{s}$

$X = 1.996 \text{ mm}$

$\sigma_H = 4.6 \text{ kbar}$

Figure 8 Wave Profile Test Records

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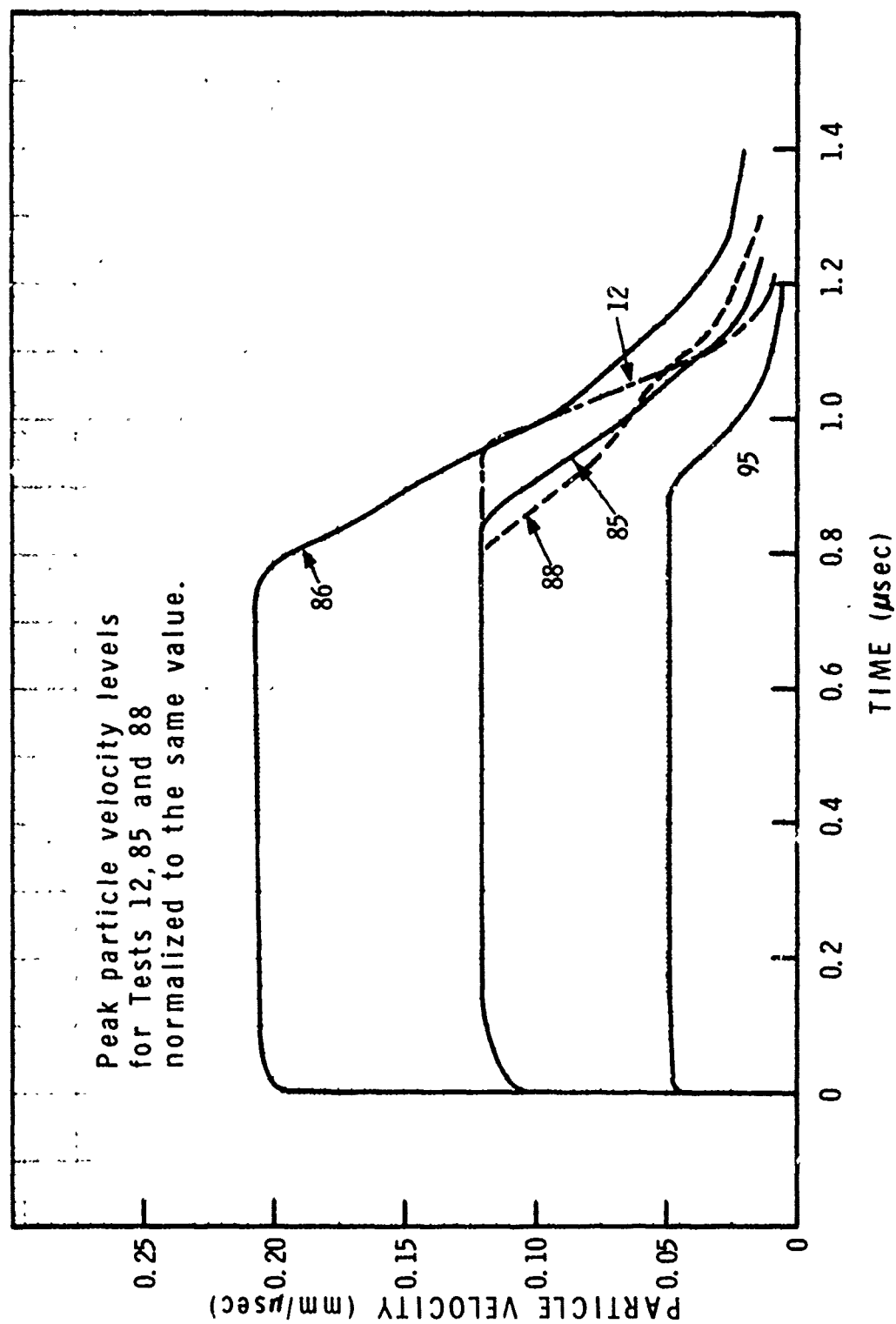


Figure 9 Wave Profiles in PMMA (Velocity Interferometer)

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TABLE III
VELOCITY INTERFEROMETER TEST DATA

Test No.	V_I Impact Velocity (mm/ μ s)	Max. Stress (kbar)	X_O Impact Thickness (mm)	X Target Thickness (mm)	X/X_O	Config.
12	0.277	5.0	1.651	1.473	0.89	PMMA \rightarrow PMMA/PMMA
85	0.246	4.4	1.542	3.391	2.20	"
86	0.407	7.5	1.527	3.414	2.24	"
88	0.250	4.5	1.590	6.459	4.06	"
89	0.256	4.6	1.463	3.269	2.23	PMMA \rightarrow PMMA
95	0.097	1.7	1.433	3.332	2.33	PMMA \rightarrow PMMA/PMMA

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increasing propagation distance is evident and there is some indication of structure in the release wave, being more pronounced at the higher particle velocity (No. 86) and at the longer propagation distance (No. 88). Similar behavior was noted by Barker.⁽²¹⁾

The magnetic foil test consisted of a multi-layered target with foil gages at the impact surface (input) and at three propagation distances. Results are shown in Figure 10, with wave arrival times calculated on the basis of a constant shock velocity of 3.07 mm/ μ sec. The initial release wave does not reduce particle velocity to zero because the impedance of the fused quartz impactor is higher than that of PMMA. The input pulse is not perfectly square due to the effects of impact tilt and gage thickness effects. As noted for the interferometer tests, there is spreading of the release wave with propagation distance although peak particle velocity attenuation has not yet occurred at the maximum distance (12.63 mm) observed.

One wave profile test was conducted on a "designed-composite." This composite consisted of aluminum rods inserted and bonded into a PMMA matrix and the test configuration is shown schematically in Figure 11. The transmitted compressive wave was measured with a 12.7 mm thick quartz gage and the test result is given in Figure 12 in terms of quartz stress. The initial portion of the wave is due to propagation down the aluminum rods and the slight structure evident in the first 0.8 μ sec is probably due to elastic-plastic effects in aluminum. Arrival of the wave propagating in the matrix gives a final stress jump to ~ 11 kbar. Wave transit times in the rods and the matrix corresponded to the elastic wave velocity in aluminum and shock wave velocity in PMMA, respectively.

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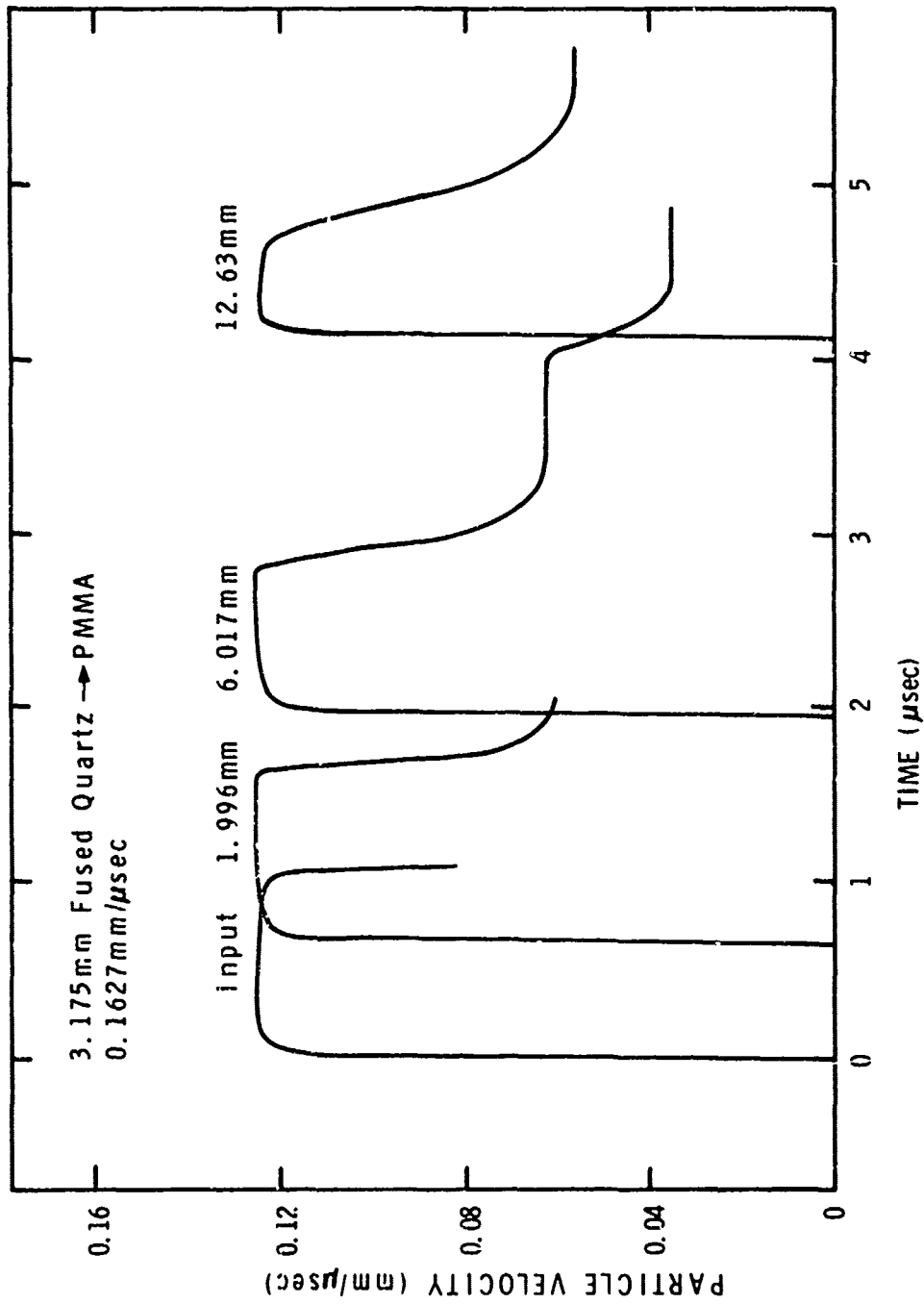


Figure 10 Wave Profiles in PMMA (Magnetic Foil Gages)

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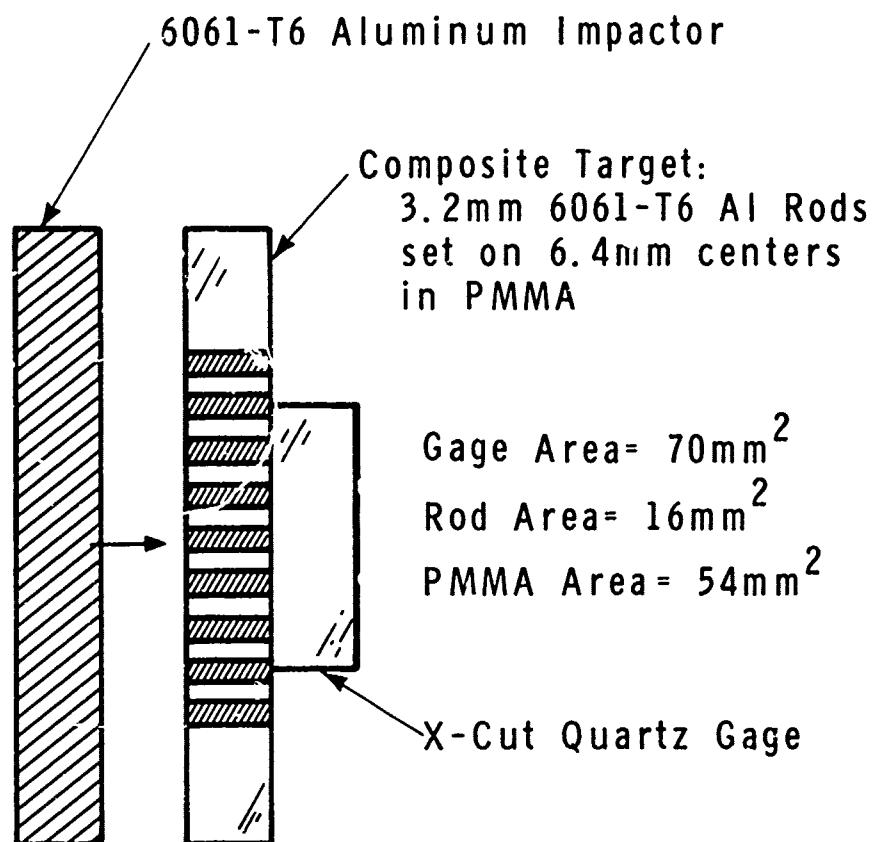


Figure 11 Designed Composite, Test Configuration

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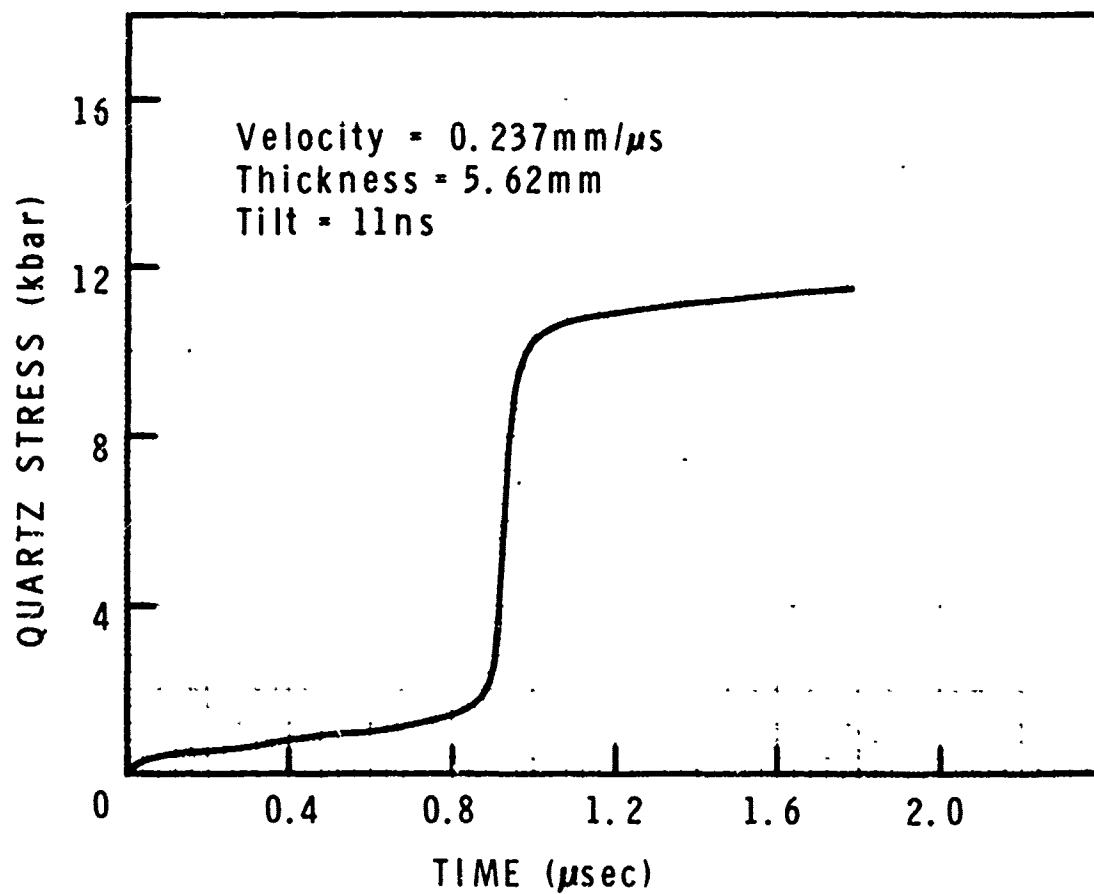


Figure 12 Compressive Wave in Designed Composite

SECTION V

SPALL FRACTURE

RECOVERY TESTS

Spall behavior of PMMA was studied by carrying out a series of impact and recovery tests.* Unlike metals, the recovered specimens did not require sectioning and metallographic preparation to determine the presence of fractures. Spall damage could be observed simply by viewing the specimen in transmitted light parallel to the shock wave propagation direction. However, the onset of spall in PMMA was evidenced by the appearance of one or a few "penny-like" cracks several millimeters in diameter, rather than numerous small voids or cracks. This, together with poor reproducibility apparently due to material effects, made determination of an incipient spall level even more difficult than with metals.

Incipient spall results for three times of loading (i.e., three impactor thicknesses), are given in Figure 13, along with earlier spall data of Charest⁽³²⁾ and Keller⁽³³⁾. The present data show a dependence of incipient spall velocity on time of loading. This is in agreement with the results of Charest who also used a gun-launched flat-plate technique. Keller showed a time-independent behavior, however, his data are for an exploding-foil technique. All three sets of data show different incipient spall levels, which is probably due to a number of factors, the most important of which are material variability, influence of differences in the X/X_0 ratio, and different incipient spall criteria.

* The fracture of PMMA under uniaxial stress conditions has been studied in detail by Beardmore and Johnstone.^(30,31)

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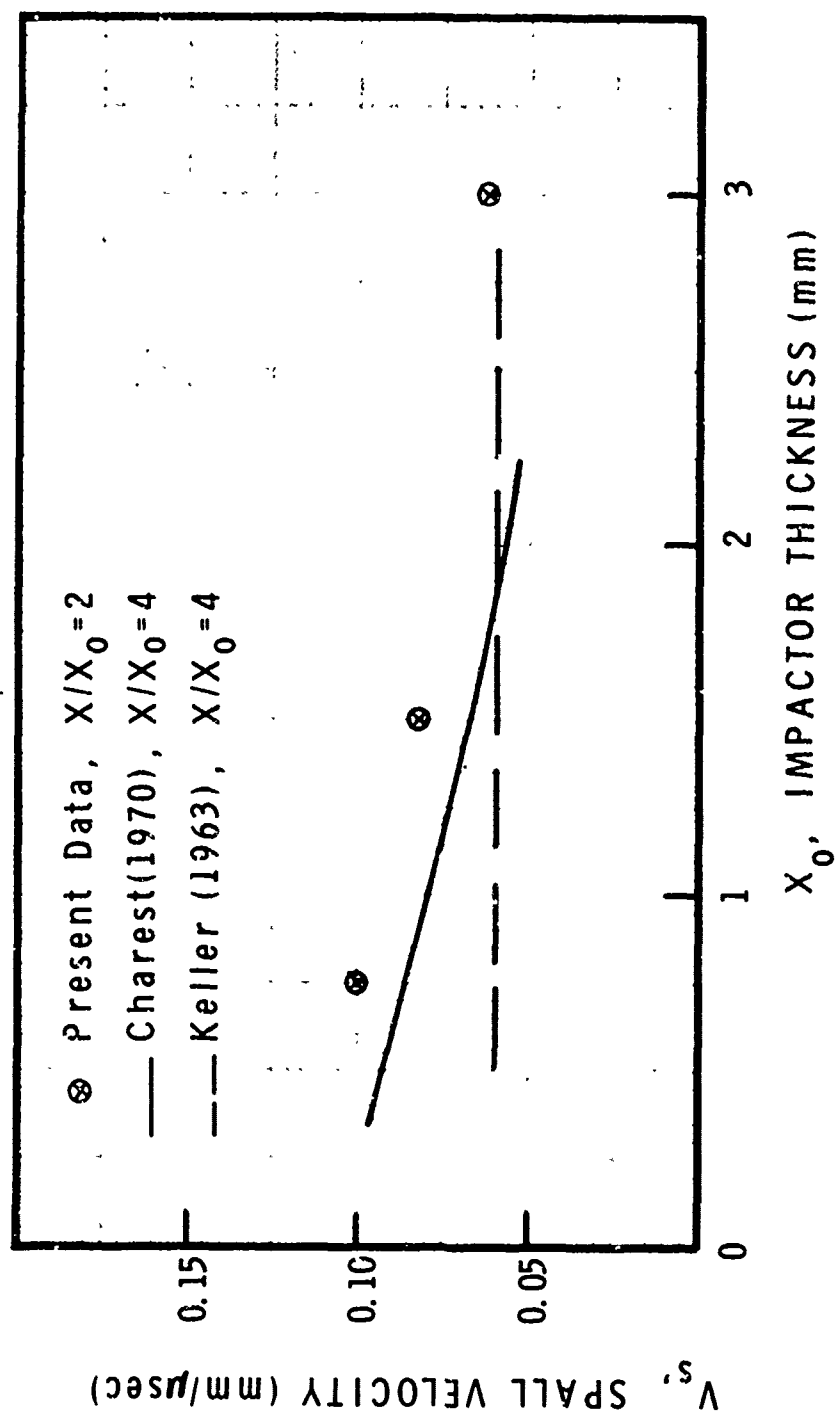


Figure 13 Spall Data for PMMA

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Fractures appearing at the incipient spall level are difficult to photograph optically because of depth of field problems and the relatively small difference in light transmission characteristics between incipient fractures and unfractured material. Optical photographs of specimens from tests $\sim 10\%$ above the incipient spall level are shown in Figure 14. (The dark "shadow" in the center of the righthand picture is a crack extending from the specimen surface to the spall plane). Although, of approximately the same total area, the fracture region has a considerably different appearance for the two specimens, the reason for which is not understood.

An optical photomicrograph of a completely spalled specimen is shown in Figure 15. The fracture surface has the appearance of a "polycrystalline" structure with extensive cracking or crazing within the "grains". The fracture subsurfaces were generally flat and normal to the shock propagation direction, although they were not all in the same plane (dispersion of ~ 0.2 mm). A scanning electron micrograph of the same specimen is shown in Figure 16.

A spall test was performed on a sample of FF-17 phenolic being studied under a related program.⁽³⁴⁾ This material is a cured phenolic resin with a density of 1.29 g/cc. The specimen was completely spalled and the resulting fracture surface is shown in Figures 17 and 18. The optical photomicrographs (Figure 17) indicate that fracture initiated at a large number of points or defects and then propagated radially from these points, giving a "sunburst" pattern. These fracture zones were typically 0.1 to 0.4 mm diameter and, like PMMA, were generally flat and normal to the shock propagation direction. The scanning microfractographs show in more detail the development of this pattern around the central initiation point.

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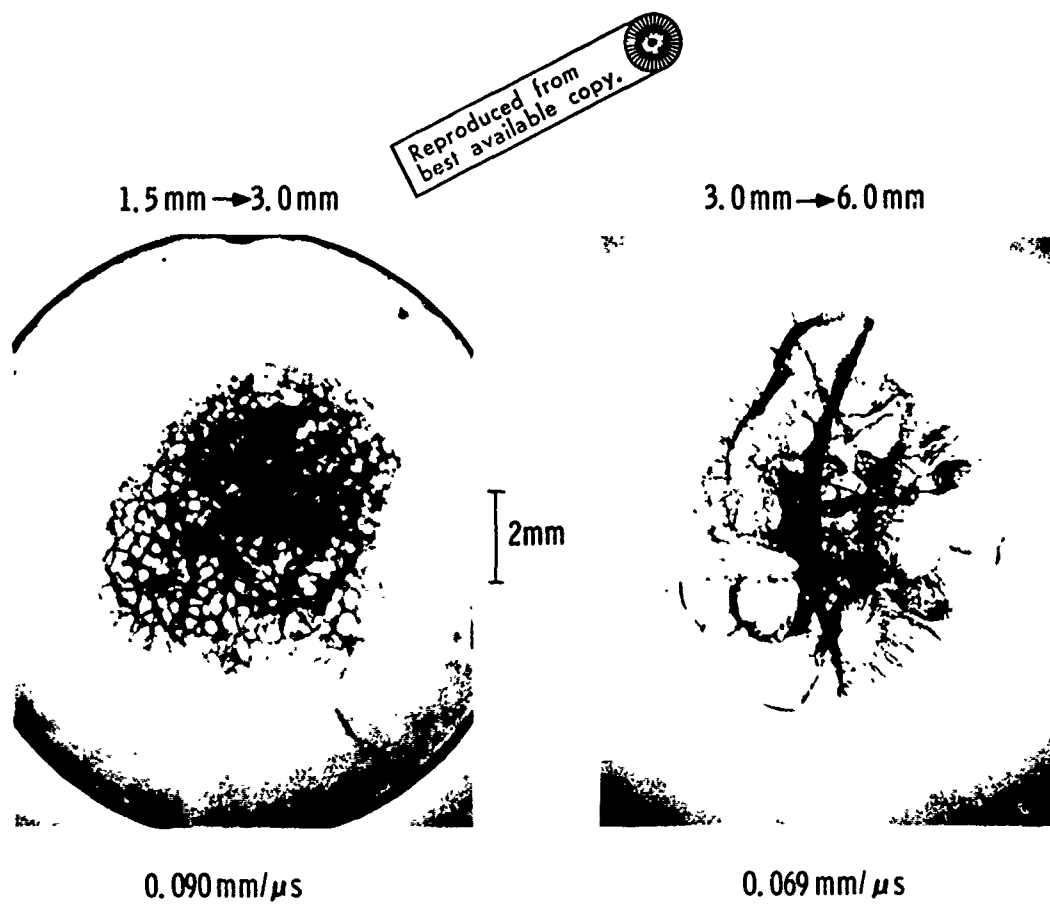


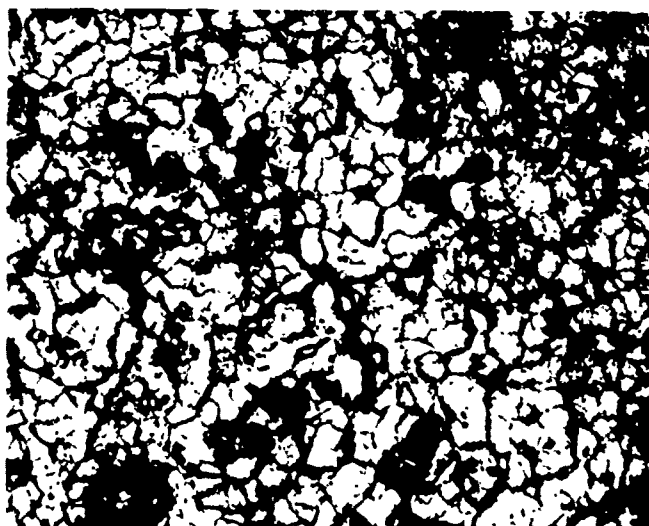
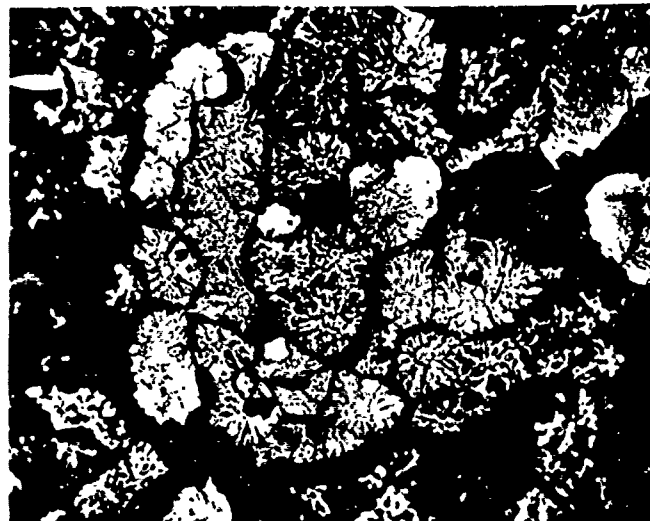
Figure 14 Spall Fractures in PMMA

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Pictures taken with transmitted light normal
to spall plane with $0.4\mu\text{m}$ sputtered gold film

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best available copy.

0.25mm



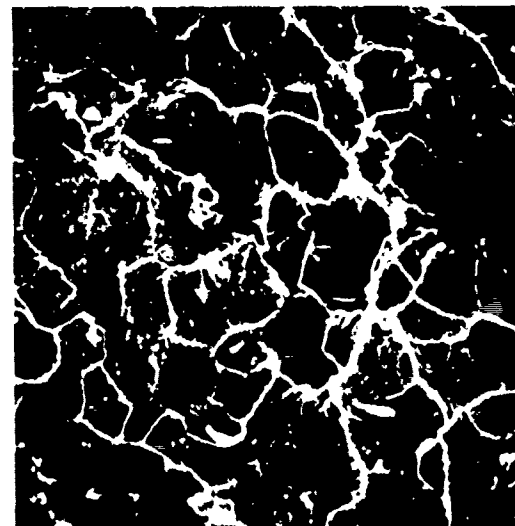
1mm

Figure 15 Spall Fracture Surface in PMMA (1.5 mm + 3.00 mm, 0.157 mm/ μsec)

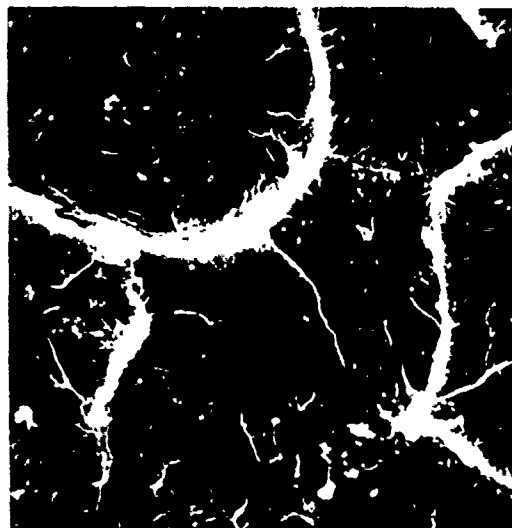
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SPUTTERED 0.4 μ m GOLD SURFACE

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200 μ m



30 μ m

Figure 16 Spall Fracture Surface in PMMA, Scanning Microfractograph, 0° (1.5 mm \rightarrow 3.0 mm, 0.157 mm/ μ sec)

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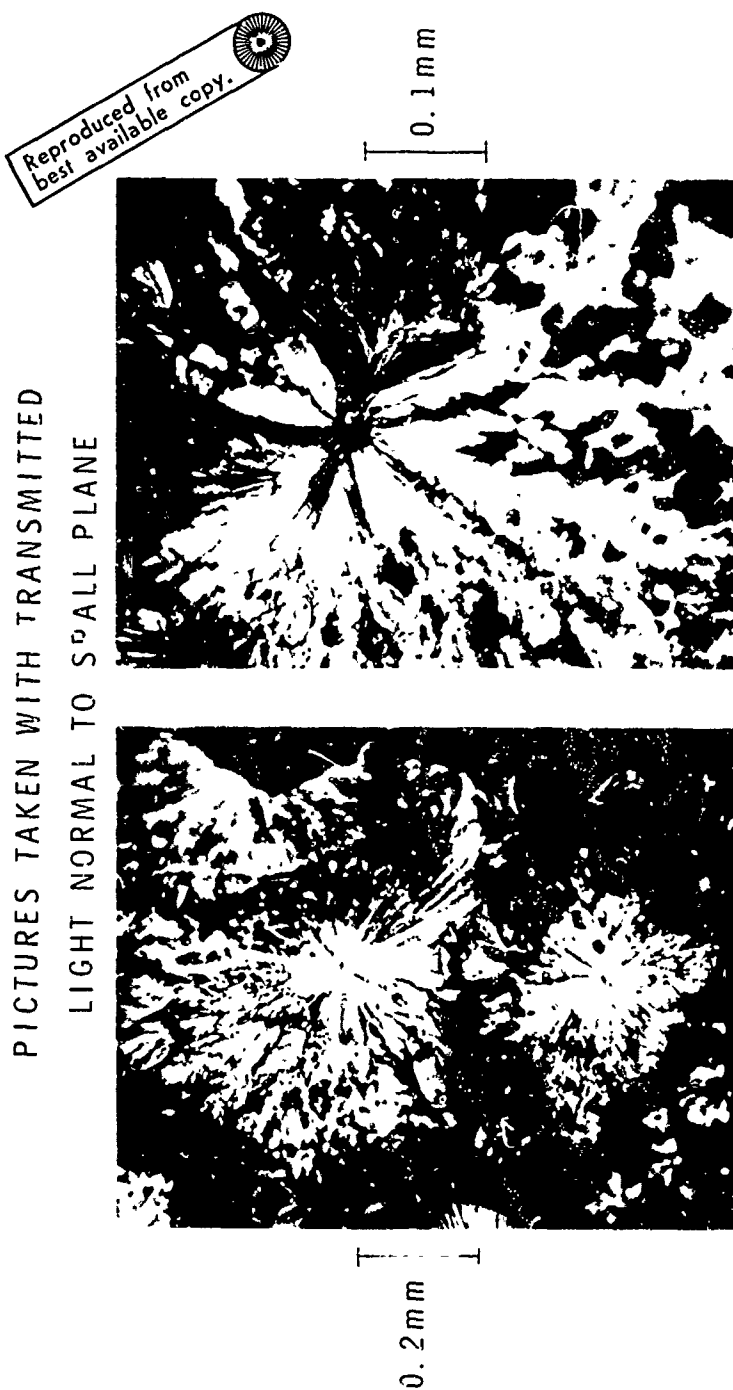


Figure 17 Spall Fracture Surfaces in FF-17 Phenolic

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SPUTTERED 0.1 μ m RHODIUM SURFACE

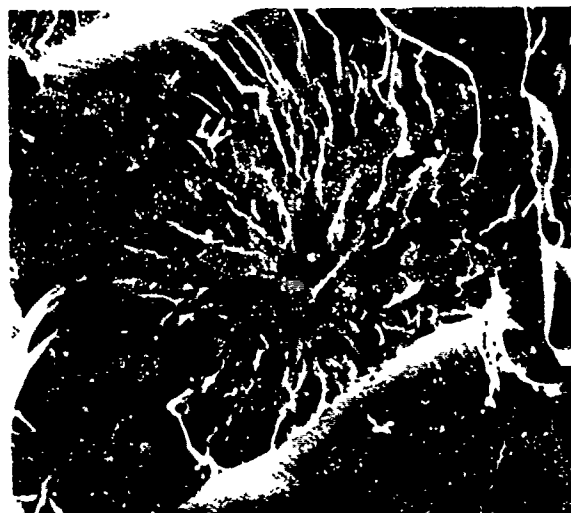


Figure 18 Spall Fracture Surface in FF-17 Phenolic,
Scanning Microfractograph, 45°

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SPALL WAVE PROFILE

One test was conducted to obtain a complete wave profile under spall-producing conditions and the result is shown in Figure 19 (see Table III for test conditions). A good, noise-free velocity interferometer record was obtained, however, analysis of this record was difficult in that the expected wave profile is not entirely consistent with the apparent fringe pattern.* The compressive wave and release wave arrival are clear and agree with the wave profiles discussed earlier. The fringe pattern does not give an obvious fringe reversal after release wave arrival and therefore several interpretations of the release path are possible. The solid line shown implies a reversal at point A, which gives a pullback or release wave reversal⁽³⁶⁾ of ~ 0.07 mm/ μ sec. This pullback is related to the spall strength of the material, and for metals⁽²⁻⁵⁾ it is always less than the incipient spall velocity for a given set of impact conditions. For the conditions shown, the incipient spall velocity in PMMA is ~ 0.08 mm/ μ sec, which supports assumption of a reversal point at A. The dashed line is for no reversal at A. Although there is a fringe beginning at point B, the absence of a reversal between A and B would give a wave profile beyond B that is not reasonable. The solid wave would continue up above the initial peak velocity level while the dashed curve would continue down towards zero velocity, neither of which is consistent with observed spall profile behavior for other materials.

* Barker gives a detailed discussion of fringe analysis in Reference 20 and 35.

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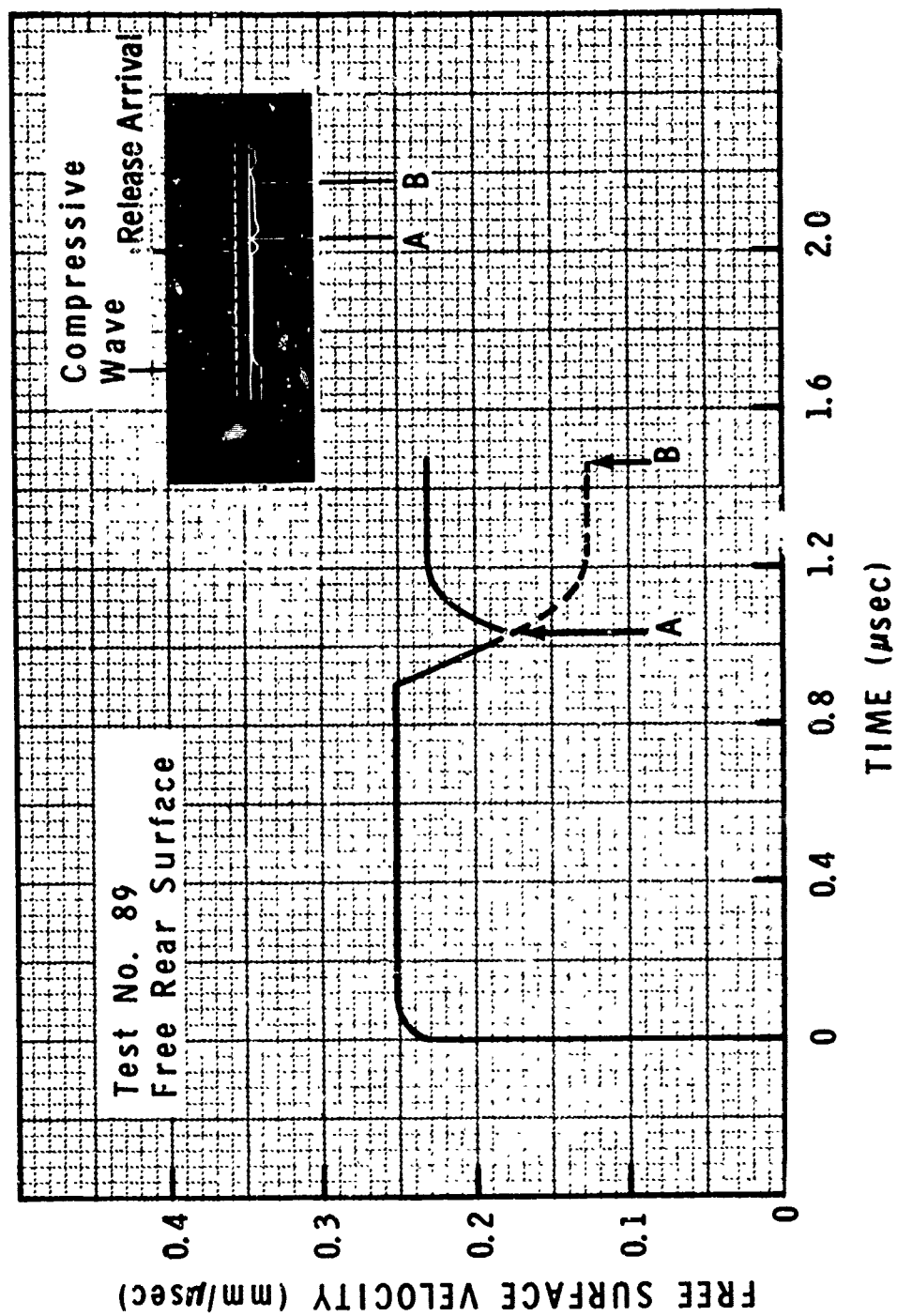


Figure 19 Spall Wave Profile in PMMA

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